# SCHOENFLIES SOLUTIONS OF CONFORMAL BOUNDARY VALUES MAY FAIL TO BE SOBOLEV 

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#### Abstract

There exists a planar Jordan domains $\Omega$ with 1-Hausdorff dimensional boundary such that, for any conformal map $\varphi: \mathbf{D} \rightarrow \Omega$, any homeomorphic extensions to the entire plane of either $\varphi$ or $\varphi^{-1}$ cannot be in $W_{\text {loc }}^{1,1}$ class (or even not in $B V_{\text {loc }}$ ).


## 1. Introduction

Let $\Gamma \subset \mathbf{C}$ be a Jordan curve, namely there exists a homeomorphism $\phi: \mathbf{S}^{1} \rightarrow \Gamma$, where $\mathbf{C}$ is the complex plane and $\mathbf{S}^{1}$ denotes the boundary of the unit disk $\mathbf{D}$. According to Jordan curve theorem, the curve $\Gamma$ divides $\mathbf{C}$ into two components, and we call the bounded component a Jordan domain.

Jordan-Schoenflies theorem states that any homeomorphism between two Jordan curves on $\mathbf{C}$ can be extended to a homeomorphism between the entire $\mathbf{C}$; see [12, Corollary 2.9]. To be more precise, given two Jordan domains $\Omega_{1}$ and $\Omega_{2}$ and a homeomorphism $\varphi: \partial \Omega_{1} \rightarrow \partial \Omega_{2}$, there exists a homeomorphism $\Phi$, which we call a Schoenflies solution of the boundary value $\varphi$, from $\mathbf{C}$ to $\mathbf{C}$ such that the restriction of $\Phi$ to $\Gamma_{1}$ coincides with $\varphi$. Then a natural question arises:

Question 1.1. Given two Jordan domains $\Omega_{1}, \Omega_{2} \subset \mathbf{C}$ together with a homeomorphism $\varphi: \partial \Omega_{1} \rightarrow \partial \Omega_{2}$, what is the best regularity of Schoenflies solutions of the boundary value $\varphi$ ?

Certainly the answer to this question depends on the given boundary value and the geometry of both $\Omega_{1}$ and $\Omega_{2}$. Let us recall some known results. If $\Omega_{2}$ is bounded by a smooth Jordan curve, then by the techniques from differential topology for each conformal map we can find a smooth Schoenflies solution to any homeomorphism from $\mathbf{S}^{1}$ onto $\partial \Omega$. Assume that $\varphi: \mathbf{S}^{1} \rightarrow \partial \Omega_{2}$ is quasisymmetric, via Douady-Earle extension theorem there exists a $K$-quasiconformal Schoenflies solution $\Phi$. By [1], we further have that both $\Phi$ and $\Phi^{-1}$ are in $W_{\text {loc }}^{1, p}(\mathbf{C})$ for any $p<2 K /(K-1)$. Recently Koskela, Pankka and the author have been working on a version of this result for domains satisfying Gehring-Martio conditions [8].

Recall Carathéodory's theorem states that, given any two Jordan domain $\Omega_{1}$ and $\Omega_{2}$, every conformal map $\varphi: \Omega_{1} \rightarrow \Omega_{2}$ can be continuously extended to the boundary as a homeomorphism $\varphi: \bar{\Omega}_{1} \rightarrow \bar{\Omega}_{2}$. We abuse $\varphi$ here. In this paper we investigate Question 1.1 with the boundary value given by Carathéodory's theorem, namely a conformal boundary value.

The main result of this paper is the following.

[^0]Theorem 1.2. There exists a Jordan domain $\Omega \subset \hat{\mathbf{C}}$ with 1-Hausdorff dimensional boundary such that, any Schoenflies solution of any conformal boundary value $\varphi: \mathbf{S}^{1} \rightarrow \partial \Omega$ or $\phi: \partial \Omega \rightarrow \mathbf{S}^{1}$ is not in $W_{\text {loc }}^{1,1}(\mathbf{C})\left(\right.$ even not in $B V_{\text {loc }}(\mathbf{C})$ ).

This result indicates that, in general, one cannot expect the regularity of Schoenflies solutions to a given boundary value to be better than homeomorphism; even if the boundary value is given by a (extended) conformal map (which is a quite natural choice). Thus, geometric assumptions on the Jordan domain in question and (energy) controls on the boundary value are necessary. One can see e.g. [2, 13, 6, 9] for recent results in this direction. Especially in the very recent paper [10] Koski and Onninen give positive answers to Question 1.1 under certain circumstances.

The notation in the paper is quite standard. The Euclidean distance between two sets $A, B \subset \mathbf{R}^{2}$ is denoted by dist $(A, B)$. We denote by $\ell(\gamma)$ the length of a curve $\gamma$. For a set $A \subset \mathbf{R}^{2}$, we write its boundary as $\partial A$, and its closure as $\bar{A}$, respectively, with respect to the Euclidean topology. We use the notation $\mathcal{H}^{1}$ for 1-dimensional Hausdorff measure.

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## 2. Proof of Theorem 1.2

Define the inner distance with respect to $\Omega$ between $x, y \in \Omega$ by

$$
\operatorname{dist}_{\Omega}(x, y)=\inf _{\gamma \subset \Omega} \ell(\gamma)
$$

where the infimum runs over all curves joining $x$ and $y$ in $\Omega$.
2.1. Schoenflies solution of conformal boundary value $\varphi: \mathrm{S}^{1} \rightarrow \partial \Omega$. The idea of the proof is that, we construct a Jordan domain $\Omega \subset \hat{\mathbf{C}}$ satisfying that there exists a (Cantor) set $E \subset \partial \Omega$ such that,
(i) for any conformal $\varphi: \mathbf{D} \rightarrow \Omega$, i.e. for any conformal boundary value, we have

$$
\mathcal{H}^{1}\left(\varphi^{-1}(E)\right)>0 ;
$$

(ii) for any point $x$ in the complementary domain $\widetilde{\Omega}$,

$$
\operatorname{dist}_{\widetilde{\Omega}}(x, E \backslash\{(0,0),(1,0)\})=\infty .
$$

If such a Jordan domain exists (see Lemma 2.1 below), then by (i) and (ii), any Schoenflies solution of the conformal boundary value $\varphi$ is not in $W_{\text {loc }}^{1,1}$ (even not in $B V_{\text {loc }}$ ) by Fubini's theorem; indeed, such a solution maps a family of radial segments in the exterior of the unit disk (with finite length) into a family curves of infinite length in $\widetilde{\Omega}$. By calculating in the polar coordinate we know that such a map cannot be in $W_{\text {loc }}^{1,1}$ (even not in $B V_{\text {loc }}$ ). Hence Theorem 1.2 follows.

We first construct a Jordan curve $\Gamma$ in the plane. Towards this, let us recall the construction of a fat Cantor set $E \subset[0,1]$ on the real axis. Let $C_{0}=I_{0,1}=[0,1]$ and $C_{i}$ with $i \geq 1$ recursively as follows: When $I_{i, j}=[a, b]$ has been defined, let

$$
I_{i+1,2 j-1}=\left[a, \frac{a+b-4^{-i}}{2}\right] \quad \text { and } \quad I_{i+1,2 j}=\left[\frac{a+b+4^{-i}}{2}, b\right] ;
$$

i.e. we remove an open interval of length $4^{-i}$ from the middle of the interval $I_{i, j}$. Then we set

$$
C_{i}=\bigcup_{j=1}^{2^{i}} I_{i, j} .
$$

The set $E$ is finally given by

$$
E=\bigcap_{i=1}^{\infty} C_{i} .
$$

A simple calculation shows that, for every $i \in \mathbf{N}$ and $1 \leq j \leq 2^{i}$, each interval $I_{i, j}$ has length

$$
\begin{equation*}
\frac{2^{i}+1}{2^{2 i+1}} \in\left(2^{-i-1}, 2^{-i}\right] \tag{2.1}
\end{equation*}
$$

Thus $C_{i}$, and hence $E$ is well-defined. Moreover, $E$ has positive $\mathcal{H}^{1}$-measure; note that at step $i, i \geq 1$ there are $2^{i}$ intervals removed with total length $2^{-i-1}$.

We now construct a sequence of simple curves $\gamma_{i}$ based on the construction of $E$. Again we proceed inductively according to the index $i$. For $i \in \mathbf{N}$ and $1 \leq j \leq 2^{i}$, denote by $I_{i, j}^{\prime} \subset I_{i, j}$ the interval removed from $I_{i, j}$ in the construction of $E$. Let $\gamma_{0}$ be the interval $[0,1]$. When $\gamma_{i-1}, i \geq 1$ has been defined, we replace every open interval $I_{i, j}^{\prime}, 1 \leq j \leq 2^{i}$, contained in $\gamma_{i-1}$, by a curve

$$
\gamma_{i, j}=\partial\left(I_{i, j}^{\prime} \times\left[0,2^{-i}\right]\right) \backslash\left(I_{i, j}^{\prime} \times\{0\}\right),
$$

consisting of three line segments, where $\times$ means the Cartesian product. We then obtain $\gamma_{i}$. See Figure 1.


Figure 1. The curve $\gamma_{2}$ is shown in the figure. In the previous steps the intervals $I_{0,1}^{\prime}, I_{1,1}^{\prime}, I_{1,2}^{\prime}$ were replaced by curves $\gamma_{0,1}, \gamma_{1,1}, \gamma_{1,2}$, and in the current step four more intervals are replaced.

Since $\left\{\gamma_{i}\right\}$ (under suitable parameterizations) is a Cauchy sequence of curves in the plane with respect to the supremum distance, the limit $\gamma$ exits and is a curve. Moreover, $\gamma$ is simple.

For fixed $n \in \mathbf{N}$, there are $2^{n+1}-1$ curves $\gamma_{i, j}$ intersecting $\mathbf{R} \times\left(2^{-n-1}, 2^{-n}\right]$. Indeed, if $\gamma_{i, j} \cap\left(\mathbf{R} \times\left(2^{-n-1}, 2^{-n}\right]\right) \neq \emptyset$, then $i \leq n$. The distance between any two of these curves is strictly larger than $2^{-n-1}$ by (2.1).

We next construct a sequence of new curves $\Gamma_{n}$ according to the index $n$. First of all define $\Gamma_{0}=\gamma$. When $\Gamma_{n-1}, n \geq 1$ has been defined, we modify the segments in

$$
\gamma_{i, j} \cap\left(\mathbf{R} \times\left(2^{-n-1}, 2^{-n}\right]\right), \quad 0 \leq i \leq n
$$

to obtain $\Gamma_{n}$. Recall that $\gamma_{i, j}$ replaces the interval $I_{i, j}^{\prime}$ in the construction of $\gamma_{i}$. Denote by $a_{i, j}$ and $b_{i, j}$ the end points of $I_{i, j}^{\prime}$ with $a_{i, j}<b_{i, j}$. Then for every $1 \leq i \leq n$,

$$
\gamma_{i, j} \cap\left(\mathbf{R} \times\left(2^{-n-1}, 2^{-n}\right]\right)=\left(\left\{a_{i, j}\right\} \times\left(2^{-n-1}, 2^{-n}\right]\right) \cup\left(\left\{b_{i, j}\right\} \times\left(2^{-n-1}, 2^{-n}\right]\right)
$$

and each $1 \leq k \leq 2^{n}-1$, we replace each segment

$$
\left\{a_{i, j}\right\} \times\left[2^{-n-1}+(4 k) 2^{-2 n-3}, 2^{-n-1}+(4 k+1) 2^{-2 n-3}\right]
$$

by

$$
\begin{aligned}
A_{i, j}^{n, k}:=\partial & \left(\left[a_{i, j}-2^{-n-1}, a_{i, j}\right] \times\left[2^{-n-1}+(4 k) 2^{-2 n-3}, 2^{-n-1}+(4 k+1) 2^{-2 n-3}\right]\right) \\
& \backslash\left\{a_{i, j}\right\} \times\left[2^{-n-1}+(4 k) 2^{-2 n-3}, 2^{-n-1}+(4 k+1) 2^{-2 n-3}\right]
\end{aligned}
$$

and

$$
\left\{b_{i, j}\right\} \times\left[2^{-n-1}+(4 k+2) 2^{-2 n-3}, 2^{-n-1}+(4 k+3) 2^{-2 n-3}\right]
$$

by

$$
\begin{aligned}
B_{i, j}^{n, k}:=\partial & \left(\left[b_{i, j}, b_{i, j}+2^{-n-1}\right] \times\left[2^{-n-1}+(4 k+2) 2^{-2 n-3}, 2^{-n-1}+(4 k+3) 2^{-2 n-3}\right]\right) \\
& \backslash\left\{b_{i, j}\right\} \times\left[2^{-n-1}+(4 k+2) 2^{-2 n-3}, 2^{-n-1}+(4 k+3) 2^{-2 n-3}\right] .
\end{aligned}
$$

This gives us the new curve $\Gamma_{n}$. See Figure 2.


Figure 2. The curve $\Gamma_{2}$ is shown in the figure, with the replacement of certain segments contained in $\Gamma_{1}$ by parts of boundaries of some rectangles, receptively.

Again since $\left\{\Gamma_{n}\right\}$ (under suitable parameterizations) is a Cauchy sequence of curves in the plane with respect to the supremum distance, we conclude that $\Gamma_{n}$ converges uniformly to some curve $\Gamma_{\infty}$ as $n \rightarrow \infty$. Moreover, according to our construction $\Gamma_{\infty}$ is simple. Define

$$
\Gamma=\Gamma_{\infty} \cup(\partial([0,1] \times[-1,0]) \backslash[0,1] \times\{0\}) .
$$

Since $\Gamma_{\infty}$ is simple, then also $\Gamma$ is simple, and hence Jordan as it is closed. We denote by $\Omega$ the bounded component of $\mathbf{C} \backslash \Gamma$.

Notice that $\partial \Omega$ is a countable union of rectifiable curves, even though it does not have finite length. Since the Hausdorff dimension of a countable union of sets is the supremum of the Hausdorff dimensions of the sets, see e.g. [11, Page 81, Section 5.9], we conclude that $\partial \Omega$ is a set of Hausdorff dimension 1.

Recall the Cantor set $E$. Now let us check that the Jordan domain $\Omega$ satisfies the two properties (i) and (ii). We remark that, in our construction, for any point $x$ in $\Omega$, we have

$$
\operatorname{dist}_{\Omega}(x, E)<\infty
$$

Before showing (i), we note that, property (i) is stated with respect to all conformal maps. However, since two such Riemann maps differ from each other by a Möbius transform on the unit disk, we may assume that $\varphi(0)$ is the center of the square $[0,1] \times[-1,0]$.

Recall that the harmonic measure in the unit disk is defined via the Poisson kernel, and then in any Jordan domain via the (extended) Riemann mapping. For
$E \subset \partial \Omega$, we use $\omega\left(x_{0}, E, \Omega\right)$ to designate the harmonic measure of $E$ at $x_{0}$ in $\Omega$. It is known that $\omega\left(x_{0}, E, \Omega\right)=u\left(x_{0}\right)$ where $u$ is the (unique) harmonic function in $\Omega$ whose boundary value is the characteristic function of $E$ on $\partial \Omega$. We refer to [4] for more details.

Lemma 2.1. The Jordan domain $\Omega$ constructed above satisfies properties (i) and (ii).

Proof. Towards (i), we first observe that

$$
\begin{equation*}
\omega(\varphi(0), E, \Omega) \geq \omega(\varphi(0), E, Q)>0 \tag{2.2}
\end{equation*}
$$

where $Q$ is the open square $(0,1) \times(-1,0)$. Indeed, the first estimate comes from the comparison principle of harmonic measures, while the second inequality follows from F. and M. Riesz theorem since its 1-Hausdorff measure is strickly positive.

By the conformal invariance of harmonic measure we have

$$
\omega\left(0, \varphi^{-1}(E), \mathbf{D}\right)>0
$$

According to the definition of harmonic measures in the unit disk, we conclude (i).
To show (ii), note that in our construction, any curve in the unbounded component of $\mathbf{R}^{2} \backslash \Gamma$ towards $E \backslash\{(0,0),(1,0)\}$ has length at least $\frac{1}{2}$ in the region $\mathbf{R} \times\left(2^{-n-1}, 2^{-n}\right]$ for $n$ large enough; the curve has to oscillate $2^{n}$ times and each time it goes at least $2^{-n-1}$. This implies that any curve in the unbounded component of $\mathbf{R}^{2} \backslash \Gamma$ towards $E \backslash\{(0,0),(1,0)\}$ has infinite length. Property (ii) is complete.
2.2. Schoenflies solution of conformal boundary value $\phi: \partial \Omega \rightarrow S^{1}$. Let $\phi: \Omega \rightarrow \mathbf{D}$ be a conformal map giving the conformal boundary value via Carathéodory's theorem. By composing with a suitable Möbius map, we may assume that $\phi\left(z_{0}\right)=0$, where $z_{0}$ is the center of open square $Q=(0,1) \times(-1,0)$; in the general case the constants below will further depending on the Möbius transform. We show that any homeomorphic extension of $\phi$ is not in $W_{\text {loc }}^{1,1}$.

Towards this, recall that in the construction of $\Gamma=\partial \Omega$ we attached "arms" $A_{i, j}^{n, k}$ and $B_{i, j}^{n, k}$ to every curve $\gamma_{i, j}$. We first claim that, there exists an absolute constant $c>0$ such that, for $n \geq 3$,

$$
\begin{equation*}
\operatorname{dist}\left(\phi\left(A_{i, j}^{n, k}\right), \phi\left(B_{i^{\prime}, j^{\prime}}^{n, k^{\prime}}\right) \geq c 2^{-n}\right. \tag{2.3}
\end{equation*}
$$

whenever $I_{i, j}^{\prime}, I_{i^{\prime}, j^{\prime}}^{\prime} \subset\left[\frac{5}{32}, \frac{27}{32}\right]$ and either $i \neq i^{\prime}$ or $j \neq j^{\prime}$.
Indeed, let us fix $A_{i, j}^{n, k}$ and $B_{i^{\prime}, j^{\prime}}^{n, k^{\prime}}$. According to our construction, there exists an interval $J \in\left\{I_{n+1, j}\right\}_{j=1}^{2^{n+1}}$ such that $J \subset\left[\frac{5}{32}, \frac{27}{32}\right]$ is between $I_{i, j}^{\prime}$ and $I_{i^{\prime}, j^{\prime}}^{\prime}$. Since $\phi: \partial \Omega \rightarrow \mathbf{S}^{1}$ is a homeomorphism, by the construction of $\partial \Omega$ and the geometry of the unit circle, we have that

$$
\operatorname{dist}\left(\phi\left(A_{i, j}^{n, k}\right), \phi\left(B_{i^{\prime}, j^{\prime}}^{n, k^{\prime}}\right)\right) \geq c_{1} \mathcal{H}^{1}(\phi(J \cap E))
$$

for some absolute constant $c_{1}$. Therefore it suffices to show that $\mathcal{H}^{1}(\phi(J \cap E)) \geq c_{2} 2^{-n}$ for some absolute constant $c_{2}$.

Again by the invariance of harmonic measure under conformal map and the comparison principle of harmonic measures,

$$
\omega(0, \phi(J \cap E), \mathbf{D})=\omega\left(z_{0}, J \cap E, \Omega\right) \geq \omega\left(z_{0}, J \cap E, Q\right)
$$

According to Schwarz-Christoffel formula [12, Chapter 3.1], since $J \subset\left[\frac{5}{32}, \frac{27}{32}\right]$ is away from the corner of $Q$, we have

$$
\omega\left(z_{0}, J, Q\right) \geq c_{3} 2^{-n}
$$

for some absolute constant $c_{3}$; note that the length of $J$ is $2^{-n-2}+2^{-1} 4^{-n-1}$, and $E$ is a self-similar fat Cantor set. Therefore, we conclude (2.3) via the Poisson formula in the unit disk.

Let $\Phi$ be any Schoenflies solution of the conformal boundary value $\phi$. By (2.3), the image under $\phi$ of any vertical segment joining "neighboring arms" $A_{i, j}^{n, k}$ and $B_{i^{\prime}, j^{\prime}}^{n, k}$ in the exterior of $\Omega$ has length at least $c 2^{-n}$. Moreover, when $n \geq 3$, the intersection of the projections on the real axis of the "neighboring arms" $A_{i, j}^{n, k}$ and $B_{i^{\prime}, j^{\prime}}^{n, k}$ is an interval with length not less than $2^{-n-2}$, and there are at least $4^{n}$ pairs of those "neighboring arms" contained in $\left[\frac{5}{32}, \frac{27}{32}\right] \times\left[2^{-n-1}, 2^{-n}\right]$ up to a multiplicative constant. Therefore by Fubini's theorem we conclude

$$
\int_{\left[\frac{5}{32}, \frac{27}{32}\right] \times\left[2^{-n-1}, 2^{-n}\right]}|D \Phi| d x \geq c^{\prime} 2^{-n} 2^{-n-2} 4^{n} \geq c^{\prime} 2^{-2}
$$

for some absolute constant $c^{\prime}$. Therefore, in any Euclidean neighborhood of $E \cap$ $\left[\frac{5}{32}, \frac{27}{32}\right]$ the $W_{\text {loc }}^{1,1}$-energy of $\Phi$ is infinite, and a similar argument shows that $\Phi \notin$ $B V_{\text {loc }}$. This concludes the second part of Theorem 1.2.

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