# SPLITTING-TYPE VARIATIONAL PROBLEMS WITH $x$-DEPENDENT EXPONENTS 

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Abstract. In this article we prove regularity results for locally bounded minimizers $u: \mathbf{R}^{n} \supset$ $\Omega \rightarrow \mathbf{R}^{N}$ of functionals of the type

$$
\int_{\Omega}\left[\left(1+\left|\nabla_{1} u\right|^{2}\right)^{\frac{p(x)}{2}}+\left(1+\left|\nabla_{2} u\right|^{2}\right)^{\frac{q(x)}{2}}\right] d x
$$

where $p$ and $q$ are Lipschitz-functions and $\nabla u=\left(\nabla_{1} u, \nabla_{2} u\right)$ is an arbitrary decompositon of the gradient of $u$. Related functionals are the topic of the paper [ Br 3 ], but the situation here is not covered.

## 1. Introduction

The study of regularity properties for minimizers $u: \Omega \rightarrow \mathbf{R}^{N}$ of energies

$$
\begin{equation*}
I[u, \Omega]:=\int_{\Omega} F(\nabla u) d x \tag{1.1}
\end{equation*}
$$

where $\Omega$ denotes an open set in $\mathbf{R}^{n}$ and where $F: \mathbf{R}^{n N} \rightarrow[0, \infty)$ satisfies an anisotropic growth condition, i.e.,

$$
C_{1}|Z|^{\bar{p}}-c_{1} \leq F(Z) \leq C_{2}|Z|^{\bar{q}}+c_{2}, \quad Z \in \mathbf{R}^{n N}
$$

with constants $C_{1}, C_{2}>0, c_{1}, c_{2} \geq 0$ and exponents $1<\bar{p} \leq \bar{q}<\infty$, was introduced by Marcellini (see [Ma1] and [Ma2]) and was widely investigated by many authors in the last years, see the references at the end of the paper. Starting from the research of Esposito, Leonetti and Mingione [ELM1] it is known that in general minimizers of (1.1) stay not regular if one allows an additional $x$-dependence and considers minimizers of functionals

$$
\begin{equation*}
J[u, \Omega]:=\int_{\Omega} F(\cdot, \nabla u) d x \tag{1.2}
\end{equation*}
$$

for $F: \Omega \times \mathbf{R}^{n N} \rightarrow[0, \infty)$. Already in the autonomous situation it is well-known that we have no hope for regularity for minimizers of (1.1) if $\bar{p}$ and $\bar{q}$ are too far apart (compare the counterexamples of [Gi2] and [Ho]). The best known bound is

$$
\bar{q}<\bar{p}+2,
$$

proven in [BF1] and [ELM2]. To get better results additional assumptions are necessary. Thus Fuchs and Bildhauer consider decomposable integrands which means we have

$$
F(Z)=f(\widetilde{Z})+g\left(Z_{n}\right)
$$

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for $Z=\left(Z_{1}, \ldots, Z_{n}\right)$ with $Z_{i} \in \mathbf{R}^{N}$ and $\widetilde{Z}=\left(Z_{1}, \ldots, Z_{n-1}\right)$ (note that this condition is only an example, we could consider every other decomposition of $\nabla u$ into two parts). Bildhauer, Fuchs and Zhong assume power growth conditions for the $C^{2}$ functions $f$ and $g$ with exponents $\bar{p} \leq \bar{q}$ and get a very general regularity theory in case $\bar{p} \geq 2$ (see [BF3], [BF4] and [BFZ]). In [Br2] we generalize these statements under the assumption

$$
f(\widetilde{Z})=a(|\widetilde{Z}|) \quad \text { and } \quad g\left(Z_{n}\right)=b\left(\left|Z_{n}\right|\right)
$$

where $a$ and $b$ are $N$-functions. Here the main assumptions are ( $h$ stands for $a$ or $b$ )

$$
\frac{h^{\prime}(t)}{t} \approx h^{\prime \prime}(t)
$$

and superquadratic growth of $h$. In [Br3] we extend the results for an $x$-dependence without severe restrictions. In this paper we focus our attention on the regularity properties for minimizers of functionals of the following type

$$
\begin{equation*}
\mathscr{F}[u]:=\int_{\Omega}\left[\left(1+|\widetilde{\nabla} u|^{2}\right)^{\frac{p(x)}{2}}+\left(1+\left|\partial_{n} u\right|^{2}\right)^{\frac{q(x)}{2}}\right] d x . \tag{1.3}
\end{equation*}
$$

Now the functions

$$
a(x, t):=\left(1+t^{2}\right)^{\frac{p(x)}{2}}-1 \quad \text { and } \quad b(x, t):=\left(1+t^{2}\right)^{\frac{q(x)}{2}}-1
$$

satisfy all conditions assumed in [Br3] (if $p, q \geq 2$ ) except

$$
\begin{equation*}
\left|\partial_{\gamma} h^{\prime}(x, t)\right| \leq c h^{\prime}(x, t) \text { for all }(x, t) \in \bar{\Omega} \times \mathbf{R}_{0}^{+} \tag{1.4}
\end{equation*}
$$

and all $\gamma \in\{1, \ldots, n\}$ for a constant $c \geq 0$. Note that (1.4) is the main hypothesis to handle the terms involving derivatives with respect to $x$ in $[\mathrm{Br} 3]$. The functions $a$ and $b$ in the functional above do not fulfill (1.4), here the best estimation is

$$
\begin{equation*}
\left|\partial_{\gamma} h^{\prime}(x, t)\right| \leq c(\epsilon)\left(1+t^{2}\right)^{\frac{\epsilon}{2}} h^{\prime}(x, t) \text { for all }(x, t) \in \bar{\Omega} \times \mathbf{R}_{0}^{+} \tag{1.5}
\end{equation*}
$$

for every $\epsilon>0$ with a constant $c(\epsilon)>0$.
Let us state our new result.
Theorem 1.1. Let $u \in L_{\text {loc }}^{\infty}\left(\Omega, \mathbf{R}^{N}\right)$ be a local minimizer of (1.3) in the class $W_{\mathrm{loc}}^{1,2}\left(\Omega, \mathbf{R}^{N}\right)$ and $p, q \in W_{\mathrm{loc}}^{1, \infty}(\Omega,[2, \infty))$. Then we have
(a) partial $C^{1, \alpha}$-regularity if $p \leq q<p+2$ on $\Omega$ (for $n \geq 5$ we additionally need $\left.p>\|q-p\|_{\infty}(n-2) / 2\right)$;
(b) full $C^{1, \alpha}$-regularity for $n=2$;
(c) full $C^{1, \alpha}$-regularity for $N=1$ if $\|p-q\|_{\infty}<2$.

Remark 1.1. Results due to minimizers of functionals like in (1.3) are not found in literature. A similar problem is minimizing

$$
\int_{\Omega}\left(1+|\nabla w|^{2}\right)^{\frac{p(x)}{2}} d x .
$$

Regularity results are stated in [CM].
Our result is not restricted to the special integrand in (1.3). We can also consider functions $a, b: \Omega \times[0, \infty) \rightarrow[0, \infty)$ which satisfy all assumptions from [Br3] except
(A5) together with (1.5). This means the theorem above covers functionals of the type

$$
\mathscr{F}[u]:=\int_{\Omega}\left[a(\cdot|\widetilde{\nabla} u|)+b\left(\cdot,\left|\partial_{n} u\right|\right)\right] d x
$$

provided we have for $h \in\{a, b\}$ and all $x \in \Omega$ :
(A1) $h(x, \cdot)$ is a $N$-function (in the sense of [Ad]);
(A2) $\frac{h^{\prime}(x, t)}{t} \approx h^{\prime \prime}(x, t)$ for all $t \geq 0$;
(A3) $a(x, t) \leq c b(x, t)$ for large $t$ (can be weakened if $n=2$ or $N=1$ );
(A4) $\frac{h^{\prime}(x, t)}{t} \geq h_{0}>0$ for all $t \geq 0$;
(A5)* for every $\epsilon>0$ we can find $c(\epsilon)>0$ such that

$$
\left|\partial_{\gamma} h^{\prime}(x, t)\right| \leq c(\epsilon)\left(1+t^{2}\right)^{\frac{\epsilon}{2}} h^{\prime}(x, t)
$$

for large $t$;
(A6) $b(x, t) \leq c t^{\omega} a(x, t)$ for large $t$ for an $\omega>0$ ( $\omega$ arbitrary if $n=2$ and $\omega<2$ if $n \geq 3$ );
(A7) $\frac{h^{\prime}(x, t)}{t} \leq h^{\prime \prime}(x, t)$ for $t \geq 0$, if $\omega<1$;
(A8) $a(x, t) \geq \vartheta t^{\frac{\omega+2 \epsilon}{2}(n-2)}$ for large $t$ and $\vartheta>0$;
(A9) $\operatorname{argmin}_{y \in B} a(y, t)$ is independent of $t$ for all $B \Subset \Omega$;
(A10) $a(x, t) \leq \theta_{1} t^{\theta_{2}|x-y|} a(y, t)$ for large $t$ and all $x, y \in B\left(\theta_{1}, \theta_{2}>0\right)$.
Note that the assumptions (A7)-(A10) disappear if $n=2$ or $N=1$ and (A7) and (A8) are only important for $n \geq 5$. Further functionals which are covered by the theory in this paper but not by [Br3] are given if we define

$$
a(x, t):=\int_{0}^{t}\left(1+s^{2}\right)^{\frac{p(x)-2}{2}} s d s \quad \text { or } \quad a(x, t):=\int_{0}^{t}(1+s)^{p(x)-2} \Theta(s) d s
$$

and $b(x, t)$ replacing $p$ by $q$. Here $\Theta \in C^{1}([0, \infty),[0, \infty))$ has to perform $\Theta^{\prime}(t) t \approx \Theta(t)$, $\Theta(0)=0$ and $\Theta^{\prime}(t) \geq \Theta_{0}>0$.

Remark 1.2. Let us compare the statements of Theorem 1.1 with the power growth situation: Fuchs and Bildhauer [BF3] proved full regularity for $n=2$ in the superquadratic situation which we can exactly reproduce. In [BF4] they analyse the general vector case and get partial regularity under the assumptions $p \leq q \leq p+2$ and $q \leq p n /(n-2)$. The first one is nearly the same as in Theorem 1.1, we can not allow an equality. If we have a look at the second one this corresponds to $p>\|q-p\|_{\infty}(n-2) / 2$ in case of constants $p$ and $q$ but without equality, too. Only the scalar case is a real restriction: In [BFZ] no condition between $p$ and $q$ is needed, but we have to suppose $\|p-q\|_{\infty}<2$.

The bound $q<p+2$ for functionals with $(p, q)$-growth firstly appears in [ELM2] (in the autonomous situation).

Remark 1.3. If $n=2$ then we do not have to assume local boundedness of the minimizer. The idea to remove this is outlined in [Bi] (section 4). In 2D it is possible to consider subquadratic problems with restriction between $p$ and $q$. In this case one can follow the approach of $[\mathrm{BF} 6]$ and $[\mathrm{Br} 4]$.

From our proof follows that we do not need superquadratic growth if $N=1$. We only have to suppose $p>1$ on $\Omega$. Then the regularized problem (compare Lemma 2.2) has a Lipschitz-solution by [BF2] (Thm. 1.2).

If $n \leq 4$, then we can deduce from $p \geq 2$ and $p \leq q<p+2$ the inequality $p>\|q-p\|_{\infty}(n-2) / 2$.

## 2. Proof of Theorem 1.1

Let

$$
a(x, t):=\left(1+t^{2}\right)^{\frac{p(x)}{2}}-1 \quad \text { and } \quad b(x, t):=\left(1+t^{2}\right)^{\frac{q(x)}{2}}-1 .
$$

It is easy to prove that these functions satisfy the assumptions (A1)-(A4) from [Br3] as well as (A9) and (A10) (compare the list in the introduction). Hence we define the regularization as there (originally it was introduced in [BF6]): for $h \in\{a, b\}$ let

$$
h_{M}(x, t):=\int_{0}^{t} s g_{M}(x, s) d s
$$

where $M \gg 1$ and

$$
g_{M}(x, t):=g(x, 0)+\int_{0}^{t} \eta(s) g^{\prime}(x, s) d s, \quad g(x, t):=\frac{h^{\prime}(x, t)}{t} .
$$

Here $\eta \in C^{1}([0, \infty))$ denotes a cut-off function with the properties $0 \leq \eta \leq 1, \eta^{\prime} \leq 0$, $\left|\eta^{\prime}\right| \leq c / M, \eta \equiv 1$ on $[0,3 M / 2]$ and $\eta \equiv 0$ on $[2 M, \infty)$. From [ Br 3$]$ we quote the following properties of $h_{M}$.

Lemma 2.1. For the sequence $\left(h_{M}\right)$ we have:
(i) $h_{M} \in C^{2}(\Omega \times[0, \infty))$ is a $N$-function, $\frac{h_{M}^{\prime}(x, t)}{t} \geq h_{0}>0$ for all $x \in \Omega$, all $t \geq 0$ and uniformly in $M$;
(ii) $h_{M} \leq h$ and $h_{M}^{\prime \prime} \leq c(M)$ on $\Omega \times \mathbf{R}_{0}^{+}$;
(iii) we have for positive constants $\bar{\epsilon}, \bar{h}$

$$
\bar{\epsilon} \frac{h_{M}^{\prime}(x, t)}{t} \leq h_{M}^{\prime \prime}(x, t) \leq \bar{h} \frac{h_{M}^{\prime}(x, t)}{t}
$$

uniformly in $M$;
(iv) if we have $p \leq q$, then

$$
a_{M}(x, t) \leq \bar{c} b_{M}(x, t) \text { for all } x \in \Omega \text { and all } t \geq 0
$$

(v) (1.5) extends to $h_{M}$ uniformly in $M$ :

$$
\left|\partial_{\gamma} h_{M}^{\prime}(x, t)\right| \leq c(\epsilon)\left(1+t^{2}\right)^{\frac{\epsilon}{2}} h_{M}^{\prime}(x, t) \text { for all }(x, t) \in \Omega \times \mathbf{R}_{0}^{+}
$$

and all $\gamma \in\{1, \ldots, n\}$;
(vi) from $q-p \leq \omega$ for a positive number $\omega$ follows

$$
b_{M}(x, t) \leq c t^{\omega} a_{M}(x, t) \text { uniformly in } M
$$

(vii) $h_{M}$ and $h_{M}^{-1}$ satisfy uniform $\Delta_{2}$-conditions, which follows from (iii);
(viii) we get from (iii) and monotonicity of $h_{M}^{\prime}$

$$
\lambda h_{M}^{\prime}(x, t) t \leq h_{M}(x, t) \leq h_{M}^{\prime}(x, t) t \text { uniformly in } M .
$$

Only (v) is not the same as in [Br3] and need a slight comment: the estimation follows from

$$
\frac{h_{M}^{\prime}(x, t)}{t}=\eta(t) \frac{h^{\prime}(x, t)}{t}+\int_{0}^{t}\left\{-\frac{\eta^{\prime}(s)}{s}\right\} h^{\prime}(x, s) d s
$$

and (1.5) using $\eta^{\prime}(s) \leq 0$.

Letting $F_{M}(x, Z):=a_{M}(x,|\widetilde{Z}|)+b_{M}\left(x,\left|Z_{n}\right|\right)$ we define for $B \Subset \Omega$

$$
\mathscr{F}_{M}[w]:=\int_{B} F_{M}(\cdot, \nabla w) d x
$$

and $u_{M}$ as the unique minimizer of $\mathscr{F}_{M}$ in $u+W_{0}^{1,2}\left(B, \mathbf{R}^{N}\right)$. We obtain for $u_{M}$ the following regularity properties.

Lemma 2.2. (i) $u_{M}$ belongs to the space $W_{\text {loc }}^{2,2}\left(B, \mathbf{R}^{N}\right)$;
(ii) $a_{M}\left(\cdot,\left|\nabla \widetilde{u}_{M}\right|\right)\left|\widetilde{\nabla} u_{M}\right|^{2}$ and $b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\partial_{n} u_{M}\right|^{2}$ are elements of $L_{\text {loc }}^{1}(B)$;
(iii) if $n=2$ or $N=1$ then we have $u_{M} \in W_{\text {loc }}^{1, \infty}\left(B, \mathbf{R}^{N}\right)$;
(iv) for $\gamma \in\{1, \ldots, n\} \partial_{\gamma} u_{M}$ solves

$$
\int_{B} D_{P}^{2} F_{M}\left(\cdot, \nabla u_{M}\right)(\nabla w, \nabla \varphi) d x+\int_{B} \partial_{\gamma} D_{P} F_{M}\left(\cdot, \nabla u_{M}\right): \nabla \varphi d x=0
$$

for all $\varphi \in W_{0}^{1,2}\left(B, \mathbf{R}^{N}\right)$ with $\operatorname{spt}(\varphi) \Subset B$;
(v) $u_{M}$ is in $W^{1,2}\left(B, \mathbf{R}^{N}\right)$ uniformly bounded and we have

$$
\sup _{M} \int_{B} F_{M}\left(\cdot, \nabla u_{M}\right) d x<\infty ;
$$

(vi) if we have $u \in L_{\text {loc }}^{\infty}\left(\Omega, \mathbf{R}^{N}\right)$, then $\sup _{M}\left\|u_{M}\right\|_{\infty}<\infty$.

Proof. By construction of $F_{M}$ we obtain the following growth conditions (compare Lemma 2.2)

$$
\begin{aligned}
& \lambda|X|^{2} \leq D_{P}^{2} F_{M}(x, Z)(X, X) \leq \Lambda_{M}\left(1+|Z|^{2}\right)^{\frac{\epsilon}{2}}|X|^{2} \\
& \left|\partial_{\gamma} D_{P} F_{M}(x, Z)\right| \leq \Lambda_{M}\left(1+|Z|^{2}\right)^{\frac{1+\epsilon}{2}}
\end{aligned}
$$

for all $X, Z \in \mathbf{R}^{n N}$, all $\gamma \in\{1, \ldots, n\}$ and all $x \in \bar{B}$ for positive constants $\lambda, \Lambda_{M}$. If we follow the approach of [BF2] (Lemma 2.8 with $\alpha=0$ ) for $p=2$ and $q=$ $2+\epsilon$, we see $\nabla u_{M} \in L_{\mathrm{loc}}^{4}\left(B, \mathbf{R}^{n N}\right)$. Note that in case $\alpha=0$ modulus dependence is not necessary. From the same proof we deduce $u_{M} \in W_{\text {loc }}^{2,2}\left(B, \mathbf{R}^{N}\right)$ and so the first two statements of the Lemma. If we quote [BF2] (Thm. 1.1), then follows $u_{M} \in W_{\text {loc }}^{1, \infty}\left(B, \mathbf{R}^{N}\right)$ for $n=2$ or $N=1$ (we can choose $\epsilon$ small enough to reach $q<p(n+1) / n)$. By approximation we get (iv), which is of course valid for $\varphi \in$ $C_{0}^{\infty}\left(B, \mathbf{R}^{N}\right)$. We can adopt the last two statements from [ Br 3$]$.

Partial regularity. Now we have to prove the higher integrability stated in [Br3] (Theorem 1.1). This means we have to show

$$
\begin{equation*}
a_{M}\left(\cdot,\left|\nabla \widetilde{u}_{M}\right|\right)\left|\widetilde{\nabla} u_{M}\right|^{2}, b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\partial_{n} u_{M}\right|^{2} \in L_{\mathrm{loc}}^{1}(B) \text { uniformly. } \tag{2.1}
\end{equation*}
$$

If we follow the lines of [Br3] (Section 2), we get by Young's inequality and Lemma 2.2 (v) for a suitable cut-off function $\eta \in C_{0}^{\infty}(B)$ and $k \in \mathbf{N}$ large enough

$$
\begin{align*}
\int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\partial_{n} u_{M}\right|^{2} d x & \leq c(\eta)+c(\eta) \int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right)\left|\widetilde{\nabla} u_{M}\right|^{2 \epsilon} d x \\
& \leq c(\eta, \tau)+\tau \int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right)\left|\widetilde{\nabla} u_{M}\right|^{2} d x \tag{2.2}
\end{align*}
$$

This is a consequence of an integration-by-parts argument, a Caccioppoli-type inequality, following by standard calculations from Lemma 2.2 (iv), and finally (1.5) resp. Lemma 2.1 (v) (of course we also need the uniform bounds from Lemma 2.2 (v)).

We remark that (2.2) is the analogy of inequality (2.5) in [Br3]. Whereas (2.7) of [Br3] now reads as

$$
\begin{equation*}
\int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right)\left|\widetilde{\nabla} u_{M}\right|^{2} d x \leq c(\eta)+c(\eta) \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\partial_{n} u_{M}\right|^{2} d x \tag{2.3}
\end{equation*}
$$

If we combine (2.2) and (2.3) and choose $\tau$ small enough we get (2.1) and can pass to the limit (details can be found in [Br3]).

As usual the key to the partial regularity is the following lemma.
Lemma 2.3. Assume the assumptions of Theorem 1.1 and fix $L>0$. Then there is a $C^{*}(L)$ such that for every $\tau \in(0,1 / 4)$ exists an $\kappa=\kappa(\tau, L)>0$ with the following property: If

$$
\begin{equation*}
\left|(\nabla u)_{x, r}\right| \leq L \quad \text { and } \quad E(x, r)+r^{\gamma^{*}} \leq \kappa \tag{2.4}
\end{equation*}
$$

for a ball $B_{r}(x) \Subset \Omega$, this implies

$$
\begin{equation*}
E(x, \tau r) \leq C^{*} \tau^{2}\left[E(x, r)+r^{\gamma^{*}}\right] \tag{2.5}
\end{equation*}
$$

where $\gamma^{*} \in(0,2)$ is arbitrary and $(f)_{x, r}$ denotes the mean value of a function $f$ over the ball $B_{r}(x)$.

Here we have

$$
E(x, r):=f_{B_{r}(x)}\left|\nabla u-(\nabla u)_{x, r}\right|^{2} d y+f_{B_{r}(x)} \bar{a}\left(\cdot,\left|\nabla u-(\nabla u)_{x, r}\right|\right) d y
$$

for a small radius $r$, where $\bar{a}(x, t):=a(x, t) t^{\omega+2 \epsilon}$ and $\omega:=\|p-q\|_{\infty}<2$. The $\epsilon$-term in the definition of $\bar{a}$ is the modification of the excess function in [Br3] and compensates the additional power $\epsilon$ in (1.5). Since $\epsilon>0$ is arbitrary and $\omega<2$, we can reach $\omega+2 \epsilon<2$ and the well-definedness of the excess function follows from (2.1) by the lines of [ Br 3$]$ (section 3).

Proof of Lemma 2.3. Thanks to the modification of $\bar{a}$ in the excess function we can prove Lemma 2.3 as in [Br3]. In the proof of the strong convergence of the scaled functions we need the convergence (letting $A_{K}(r):=B_{r} \cap\left[\lambda_{m}\left|\widetilde{\nabla} u_{m}\right|>K\right], K>3 L$, $\left.B_{r} \Subset B_{1}\right)$

$$
\begin{equation*}
\lambda_{m}^{-2} \int_{\bar{A}_{K}(r)} \bar{a}\left(\lambda_{m}\left|\nabla u_{m}\right|\right) d y \longrightarrow 0, \quad m \rightarrow \infty \tag{2.6}
\end{equation*}
$$

Here $u_{m}$ as a scaling of $u$ on the unit ball and $\lambda_{m}$ converge to zero, details, also for (2.6), can be found in $[\mathrm{Br} 2]$ and $[\mathrm{Br} 3]$. Since (compare [Fu], after (3.25))

$$
\lambda_{m}^{-2} \int_{\bar{A}_{K}(r)} \bar{a}\left(\lambda_{m}\left|\widetilde{\nabla} u_{m}\right|\right) d y \leq c(r)\left(\int_{A_{K}(r)}\left|\lambda_{m} \widetilde{\nabla} u_{m}\right|^{\frac{\omega+2 \epsilon}{2} n} d z\right)^{2 / n}
$$

we conclude (2.6) by

$$
a(x, t) \geq \vartheta t^{\frac{\omega+2 \epsilon}{2}(n-2)}
$$

which follows from $p>\|p-q\|_{\infty}(n-2) / 2$ for a suitable choice of $\epsilon$ and the definition of $\omega$. Now we have

$$
\lambda_{m}^{-2} \int_{\bar{A}_{K}(r)} \bar{a}\left(\lambda_{m}\left|\widetilde{\nabla} u_{m}\right|\right) d y \leq c(r)\left(\int_{A_{K}(r)} a\left(\cdot,\left|\lambda_{m} \widetilde{\nabla} u_{m}\right|\right)^{\frac{n}{n-2}} d z\right)^{2 / n}
$$

and the r.h.s. vanishes for $m \rightarrow \infty$, for details we refer again to [ Br 2$]$ and [ Br 3$]$. For

$$
\lambda_{m}^{-2} \int_{A_{K}(r)} b\left(\left|\lambda_{m} \partial_{n} u_{m}\right|\right)\left|\lambda_{m} \partial_{n} u_{m}\right|^{\omega} d y
$$

the same arguments are applicable. This finally leads to (2.6), the last missing step in the contradiction of the proof of Lemma 2.3 from which the claim of Theorem 1.1 follows by standard arguments.

Full regularity for $\boldsymbol{n}=\mathbf{2}$. In [BF6], (2.5), the authors prove an inequality of the form (sum over $\gamma \in\{1,2\}$ )

$$
\begin{align*}
& \int_{B_{r}(z)} D_{P}^{2} F_{M}\left(\cdot, \nabla u_{M}\right)\left(\partial_{\gamma} \nabla u_{M}, \partial_{\gamma} \nabla u_{M}\right) d x \\
& \leq c(\tau)(R-r)^{-\beta}+\tau \int_{B_{R}(z)}\left(a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2}+b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)^{2}\right) d x \tag{2.7}
\end{align*}
$$

Here is $B_{r}(z) \Subset B_{R}(z) \Subset B, \tau>0$ arbitrary and $\beta>0$ a suitable exponent. On account of the $x$-dependence we have additionally to the terms in [BF6] the integral

$$
-\int_{B_{R}(z)} \eta^{2} \partial_{\gamma} D_{P} F_{M}\left(\cdot, \nabla u_{M}\right): \partial_{\gamma} \nabla u_{M} d x
$$

where $\eta \in C_{0}^{\infty}\left(B_{R}(z)\right)$ is a suitable cut-off function. Using (1.5) and the splittingstructure we estimate this by

$$
\begin{aligned}
& c \int_{B_{R}(z)} \eta^{2} a_{M}^{\prime}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2}}\left|\partial_{\gamma} \partial_{1} u_{M}\right| d x \\
& +c \int_{B_{R}(z)} \eta^{2} b_{M}^{\prime}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)\left(1+\left|\partial_{2} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2}}\left|\partial_{\gamma} \partial_{2} u_{M}\right| d x .
\end{aligned}
$$

As a consequence of Young's inequality we can bound the first term by (compare Lemma 2.1 (viii))
$\tau^{\prime} \int_{B_{R}(z)} \eta^{2} \frac{a_{M}^{\prime}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)}{\left|\partial_{1} u_{M}\right|}\left|\partial_{\gamma} \partial_{1} u_{M}\right|^{2} d x+c\left(\tau^{\prime}\right) \int_{B_{R}(z)} \eta^{2} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{\epsilon} d x$.
For $\tau^{\prime} \ll 1$ one can absorb the first integral in the l.h.s. of (2.7). Here we used the inequality

$$
\frac{a_{M}^{\prime}(\cdot,|\widetilde{Z}|)}{|\widetilde{Z}|}|\widetilde{P}|^{2} \leq c D_{P}^{2} F_{M}(x, Z)(P, P)
$$

for $Z, P \in \mathbf{R}^{n N}$ (compare Lemma 2.1, part (iii)). For the second one we obtain

$$
\begin{aligned}
& c\left(\tau^{\prime}\right) \int_{B_{R}(z)} \eta^{2} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{\epsilon} d x \\
& \leq \tau^{\prime \prime} \int_{B_{R}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2} d x+c\left(\tau^{\prime \prime}\right) \int_{B_{R}(z)}\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{2 \epsilon} d x .
\end{aligned}
$$

We can handle the r.h.s. conveniently, since we may assume $\epsilon \leq 1 / 2$ and receive (compare Lemma 2.2, part (v))

$$
\int_{B_{R}(z)}\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{2 \epsilon} d x \leq c+\int_{B_{R}(z) \cap\left[\left|\partial_{1} u_{M}\right|>1\right]} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right) d x \leq c .
$$

Analogously we can incorporate the term

$$
\int_{B_{R}(z)} \eta^{2} b_{M}^{\prime}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)\left(1+\left|\partial_{2} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2}}\left|\partial_{\gamma} \partial_{2} u_{M}\right| d x
$$

hence (2.7) follows. In [BF6] we can find the inequality

$$
\begin{align*}
& \int_{B_{\rho}(z)}\left(a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2}+b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)^{2}\right) d x \\
& \leq c(R-\rho)^{-2}+c \int_{B_{r}(z)} D_{P}^{2} F_{M}\left(\cdot, \nabla u_{M}\right)\left(\partial_{\gamma} \nabla u_{M}, \partial_{\gamma} \nabla u_{M}\right) d x \tag{2.8}
\end{align*}
$$

for $\rho \in(0, R)$, and $r=(\rho+R) / 2$. In order to show this the authors of [BF6] use an additional cut-off function and Sobolev's embedding $W^{1,1} \hookrightarrow L^{2}$, valid for $n=2$, as well as the uniform growth conditions for $a_{M}$ and $b_{M}$ (compare Lemma 2.1). In our approach we obtain on the r.h.s. of this inequality additionally the term (if we estimate $\nabla_{x} a_{M}$ and $\nabla_{x} b_{M}$ using (1.5))

$$
\left[\int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2}} d x\right]^{2}+\left[\int_{B_{r}(z)} b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)\left(1+\left|\partial_{2} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2}} d x\right]^{2}
$$

We can handle both terms in a similar way and show the proceeding for the first one. By Hölder's inequality we receive the upper bound

$$
\begin{aligned}
Y_{M}:= & {\left[\int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{s \chi} d x\right]^{\frac{2}{\chi}} } \\
& \cdot\left[\int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{\frac{\chi-s \chi}{\chi-1}}\left(1+\left|\partial_{1} u_{M}\right|^{2}\right)^{\frac{\epsilon}{2} \frac{\chi}{\chi-1}} d x\right]^{2 \frac{\chi-1}{\chi}} .
\end{aligned}
$$

Here we have $s \in(0,1)$ and $\chi \in(1,2)$ such that $s \chi>1$. For the second integral $Y_{M}^{2}$ follows by Lemma 2.2 (v)

$$
Y_{M}^{2}=\int_{B_{r}(z) \cap\left[\left|\partial_{1} u_{M}\right| \leq 1\right]} \ldots+\int_{B_{r}(z) \cap\left[\left|\partial_{1} u_{M}\right|>1\right]} \ldots \leq c+\int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right) d x \leq c
$$

Note that we have for $t \geq 1$

$$
a_{M}(x, t)^{\frac{\chi-s x}{x-1}}\left(1+t^{2}\right)^{\frac{\epsilon}{2} \frac{\chi}{\chi-1}} \leq c a_{M}(x, t)
$$

for $\epsilon$ small enough, since $s \chi>1$ (remember Lemma 2.1 (i)). Now we get, using Jensen's and Young's inequality

$$
\begin{aligned}
Y_{M} & \leq c\left[\int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{s \chi} d x\right]^{\frac{2}{\chi}} \leq c \int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2 s} d x \\
& \leq \tau^{\prime \prime \prime} \int_{B_{r}(z)} a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2} d x+c\left(\tau^{\prime \prime \prime}\right)
\end{aligned}
$$

So we have to add

$$
\begin{equation*}
\tau^{\prime \prime \prime} \int_{B_{r}(z)}\left(a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2}+b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)^{2}\right) d x \tag{2.9}
\end{equation*}
$$

on the r.h.s. of (2.8). Combining (2.7)-(2.9) we have shown (for a suitable choice of $\tau$ and $\tau^{\prime \prime \prime}$ )

$$
\begin{aligned}
& \int_{B_{\rho}(z)}\left(a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2}+b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)^{2}\right) d x \\
& \leq c(R-r)^{-\beta}+\frac{1}{2} \int_{B_{R}(z)}\left(a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)^{2}+b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)^{2}\right) d x
\end{aligned}
$$

From [Gi1] (Lemma 5.1, p. 81) we obtain uniform bounds on $a_{M}\left(\cdot,\left|\partial_{1} u_{M}\right|\right)$ and $b_{M}\left(\cdot,\left|\partial_{2} u_{M}\right|\right)$ in $L_{\text {loc }}^{2}(B)$. Now we get the uniform boundedness of $u_{M}$ in $W_{\mathrm{loc}}^{2,2}\left(B, \mathbf{R}^{N}\right)$ (compare Lemma 2.1 (i) and (2.7)) and we can reproduce the proof of [BF6] for the rest, whereby the terms which appear additionally on account of (1.5) are uncritical.

Full regularity for $N=1$. In case $N=1$ it is possible to modify the $N$ function in [Br3], (4.4). Therefore we need the inequalities

$$
\begin{equation*}
b_{M}(x, t) \leq c t^{2-2 \epsilon} a_{M}(x, t) \quad \text { and } \quad a_{M}(x, t) \leq c t^{2-2 \epsilon} b_{M}(x, t) \tag{2.10}
\end{equation*}
$$

By Lemma 2.1 (vi) this follows from $\|p-q\|_{\infty}<2$ for $\epsilon \ll 1$. So we can separate the mixed integrands of the terms

$$
\int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right) \left\lvert\, \Gamma_{n, M}^{\frac{\alpha+2 \epsilon}{2}} d x \quad\right. \text { and } \quad \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right) \widetilde{\Gamma} \widetilde{M}^{\frac{\alpha+2 \epsilon}{2}} d x
$$

which occur additionally to the integrals in [ Br 3$]$. Finally we get instead of $[\mathrm{Br} 3]$, (4.6),

$$
\begin{aligned}
& \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right) \left\lvert\, \Gamma_{n, M}^{\frac{\alpha+2}{2}} d x\right. \\
& \leq c(\eta)\left[\ldots+\int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right) \left\lvert\, \Gamma_{n, M}^{\frac{\alpha+2 \epsilon}{2}} d x+\int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right) \widetilde{\Gamma_{M}} \frac{\alpha+2 \epsilon}{2} d x\right.\right]
\end{aligned}
$$

as well as an analogous inequality for $a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right) \widetilde{\Gamma}_{M}^{\frac{\alpha+2}{2}}$ instead of [Br3], (4.7). The key for this is an integrating-by-parts argument as in the vector-valued situation (compare the step: partial regularity), the growth conditions of $a_{M}$ and $b_{M}$ and a Caccioppoli-type inequality only valid if $N=1$. Since we may assume $\epsilon \leq 1 / 2$, the first integral on the r.h.s. is bounded by (using Young's inequality)

$$
\tau \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\Gamma_{n, M}^{\frac{\alpha+2}{2}} d x+c(\tau) \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\right| \Gamma_{n, M}^{\frac{\alpha}{2}} d x
$$

for an arbitrary $\tau>0$. For the second one we can argue similarly and we obtain by absorbing the $\tau$-terms

$$
\begin{aligned}
& \int_{B} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right)\left|\Gamma_{n, M}^{\frac{\alpha+2}{2}} d x+\int_{B} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right)\right| \widetilde{\Gamma}_{M}^{\frac{\alpha+2}{2}} d x \\
& \leq c(\eta)\left[\int_{\operatorname{spt}(\eta)} \eta^{2 k} b_{M}\left(\cdot,\left|\partial_{n} u_{M}\right|\right) \left\lvert\, \Gamma_{n, M}^{\frac{\alpha}{2}} d x+\int_{\operatorname{spt}(\eta)} \eta^{2 k} a_{M}\left(\cdot,\left|\widetilde{\nabla} u_{M}\right|\right) \widetilde{\Gamma}_{M}^{\frac{\alpha}{2}} d x\right.\right] .
\end{aligned}
$$

Now we can iterate as in $[\mathrm{Br} 3]$ and obtain arbitrary high integrability of $\nabla u_{M}$ uniform in $M$ (the starting point is $\alpha=0$, see Lemma 2.2, part (v)). This is enough to end up the proof as mentioned there (full regularity follows by DeGiorgi-type arguments using Stampacchia's Lemma [St], Lemma 5.1, p. 219).

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